

Gravitational Wave Astronomy With LISA

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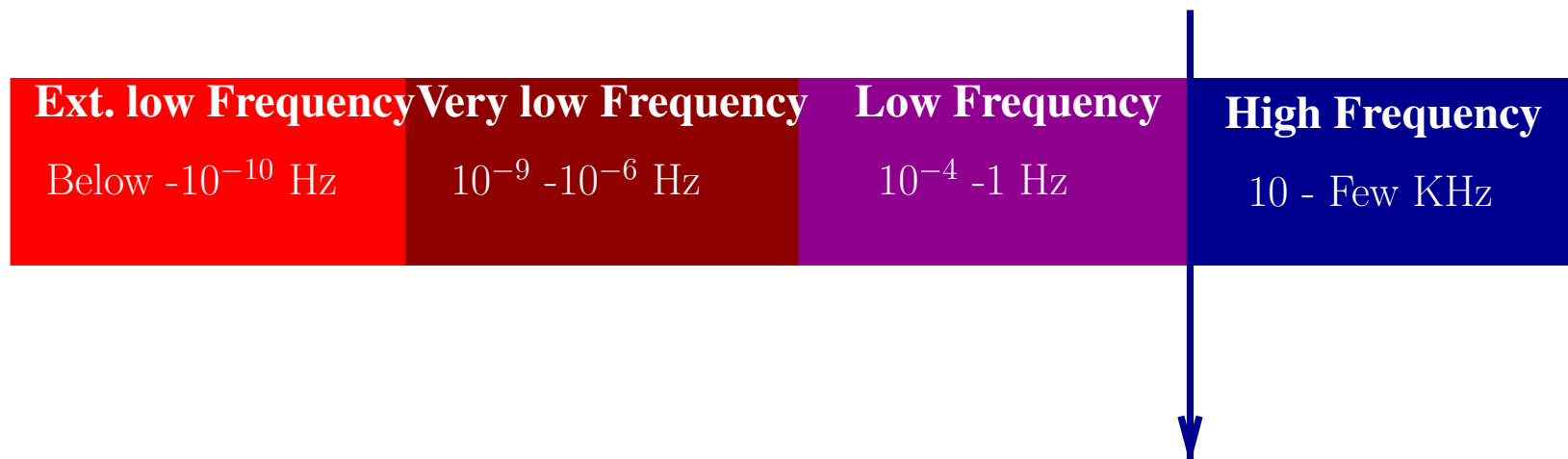
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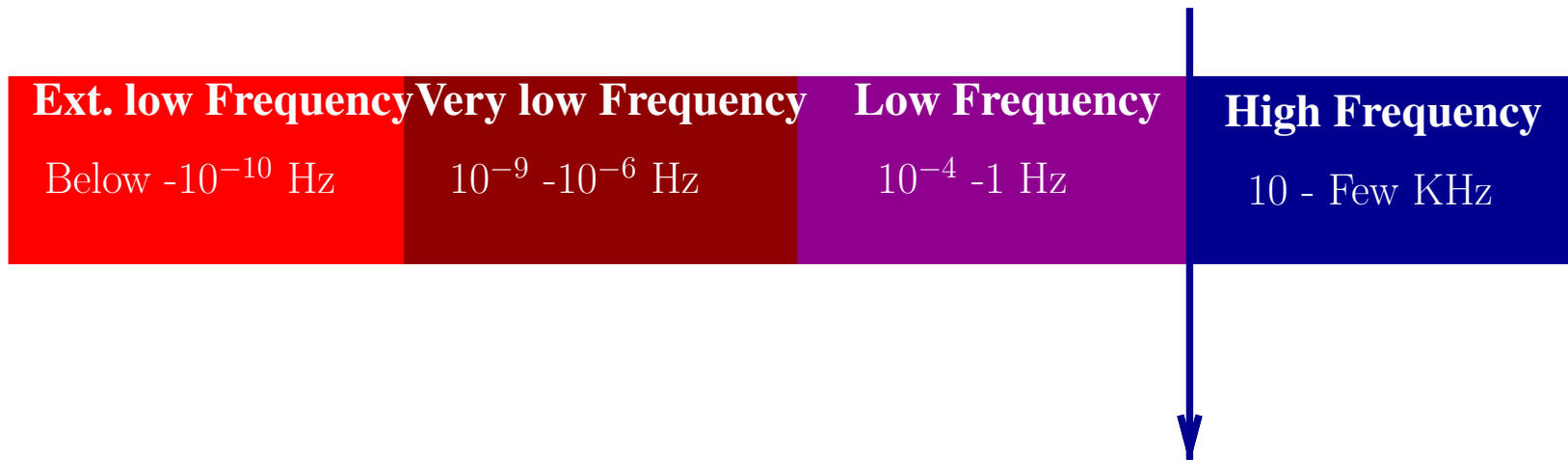


Coalescing neutron

Stellar mass Blackhole mergers

Rotating Neutron/Compact Stars

Supernova



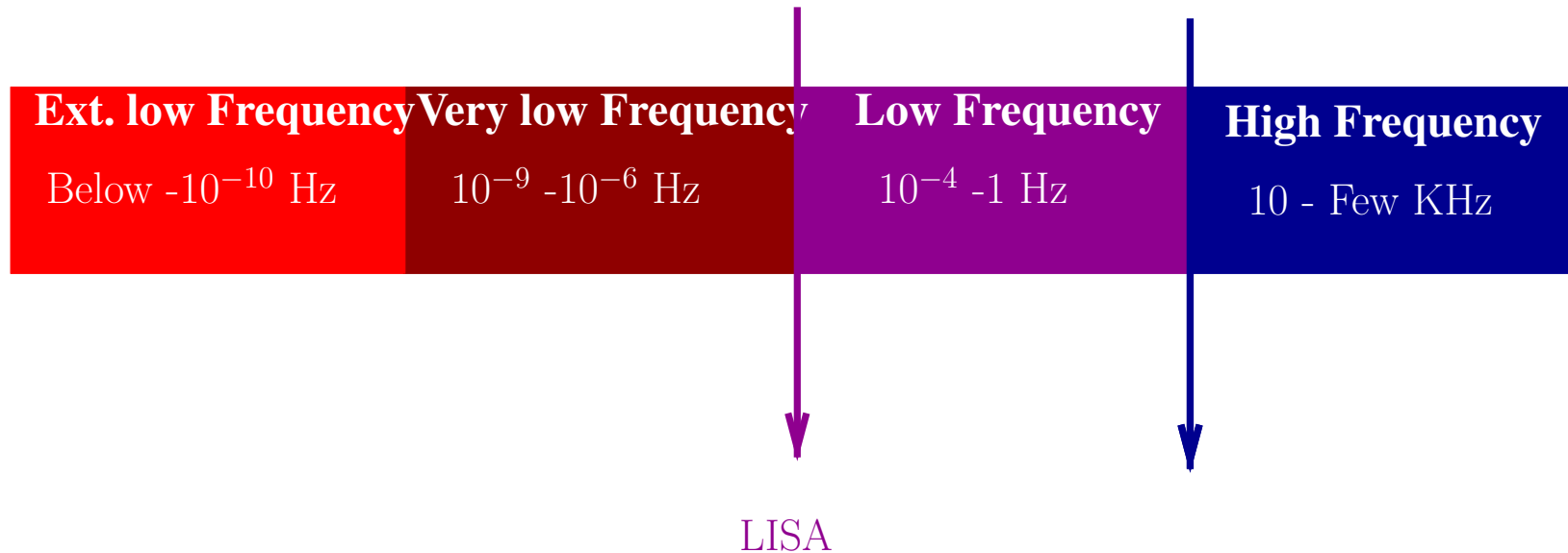
Ground Based Detectors

Such as LIGO, VIRGO etc

Coalescing SMB $< 10^7 M_s$

Galactic Binaries

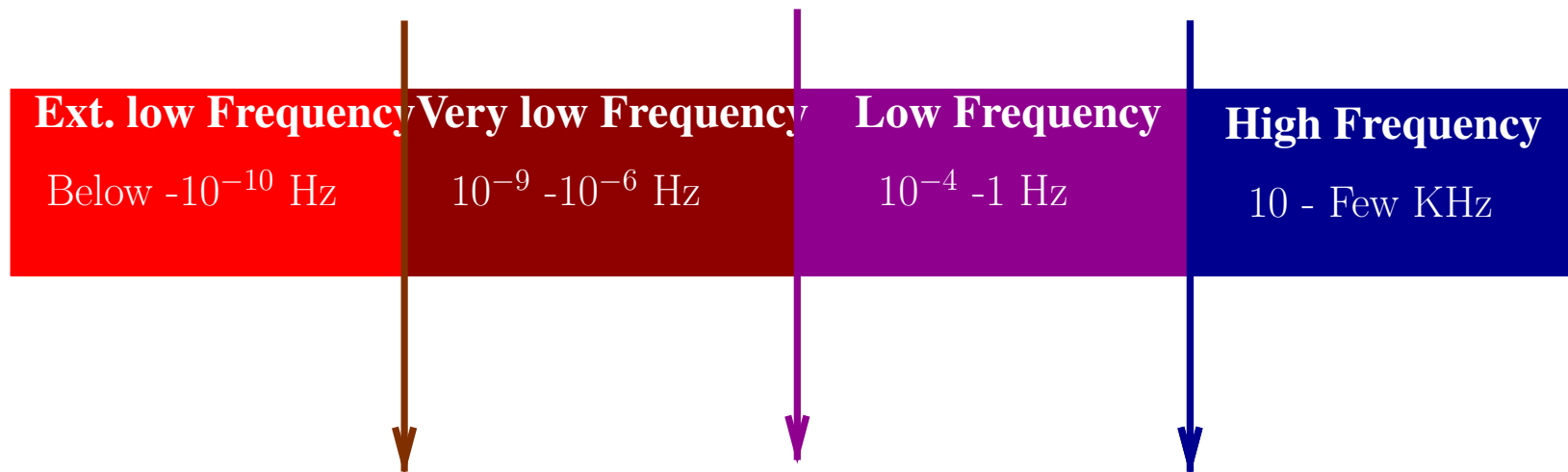
EMRI's



Coalescing SMB upto $10^{11} M_s$

Stochastic GW background

Binary stars

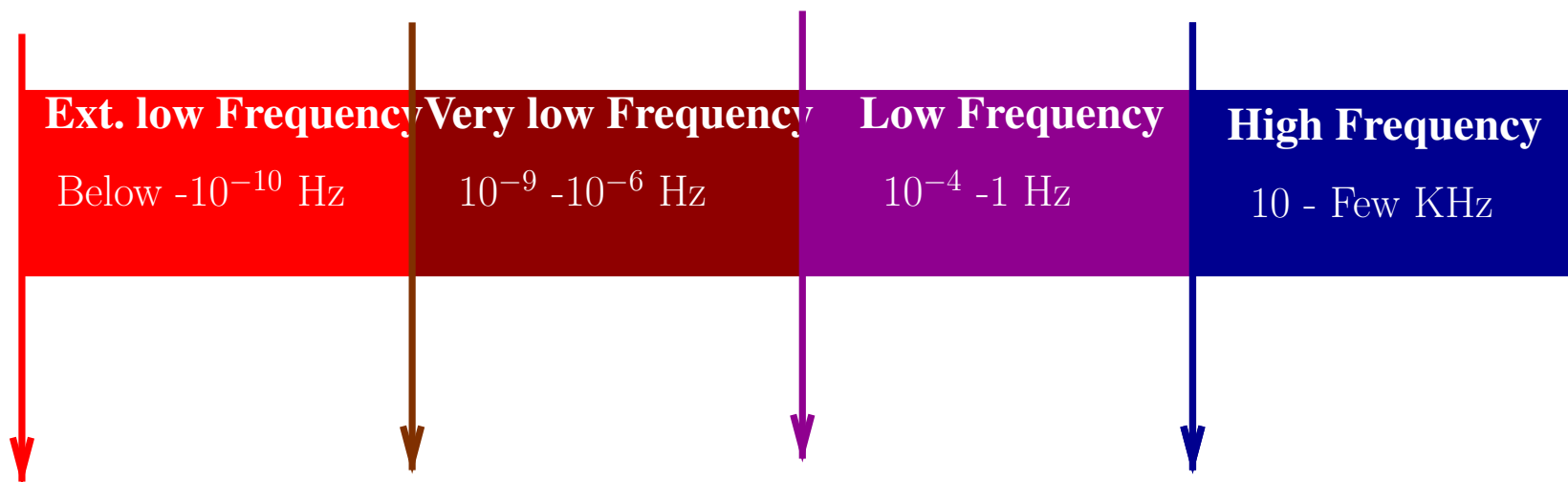


Pulsar in Binary (Taylor Pulsar)

Pulsar Timing

Doppler tracking of inter-planetary spacecrafts

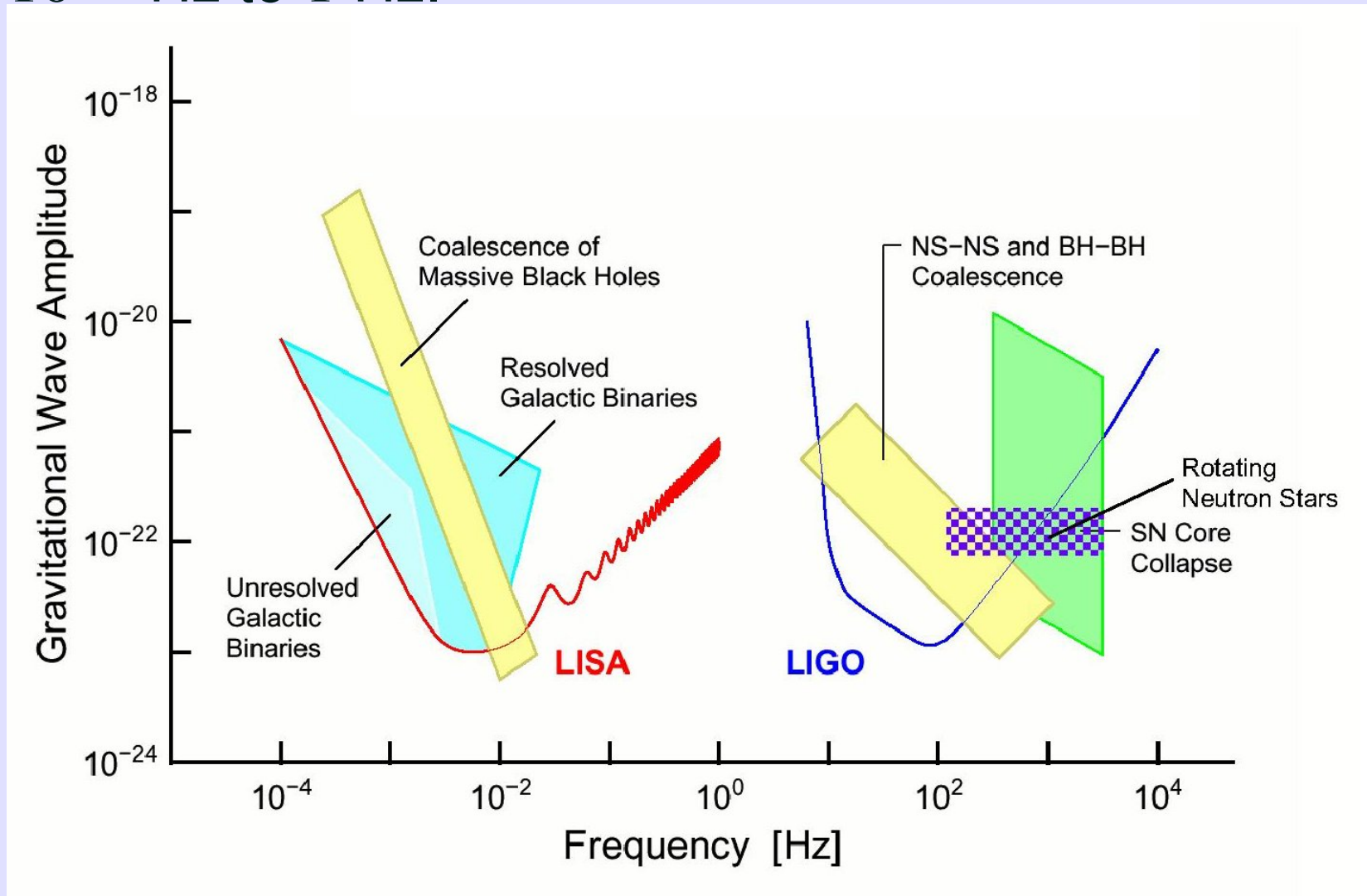
Universe



CMB

Sensitivity Curve

- ☞ LISA is sensitive to gravitational waves in the frequency range 10^{-4} Hz to 1 Hz.



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- ☞ White dwarfs have mass about $1 - 1.4 M_{\odot}$ while their radius is only 1000 KM.

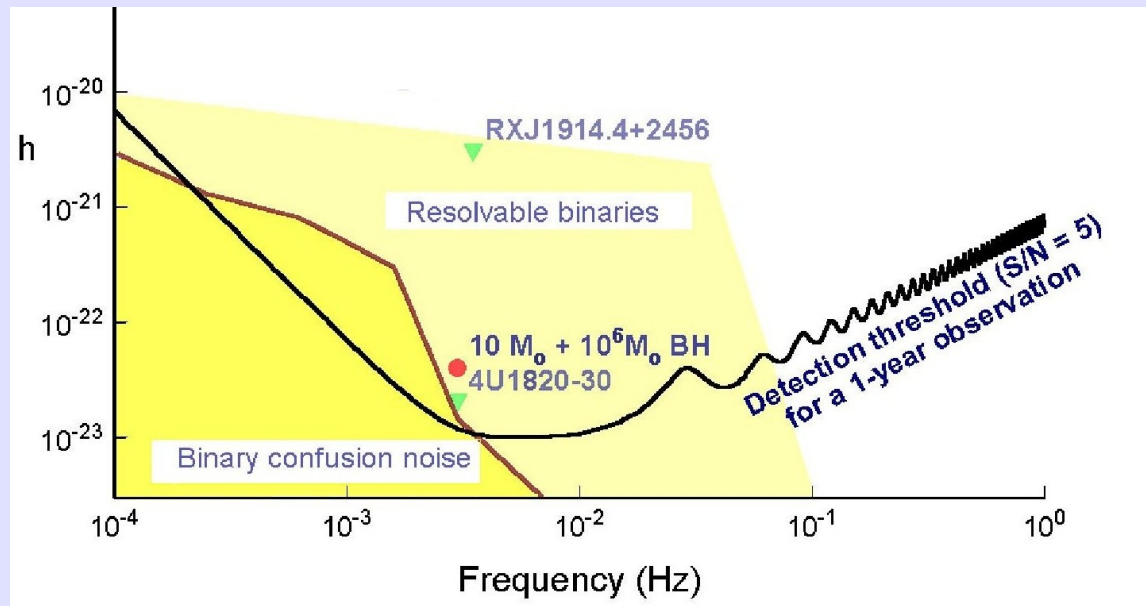
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- ☞ White dwarfs have mass about $1 - 1.4 M_{\odot}$ while their radius is only 1000 KM.
- ☞ In our Galaxy there 10^{11} stars and about 50% of them are binary system.

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Sensitivity curve for binary confusion noise

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- ☞ LISA observation can give distribution of them in our Galaxy and one can constrain the Galactic models based on LISA observations.

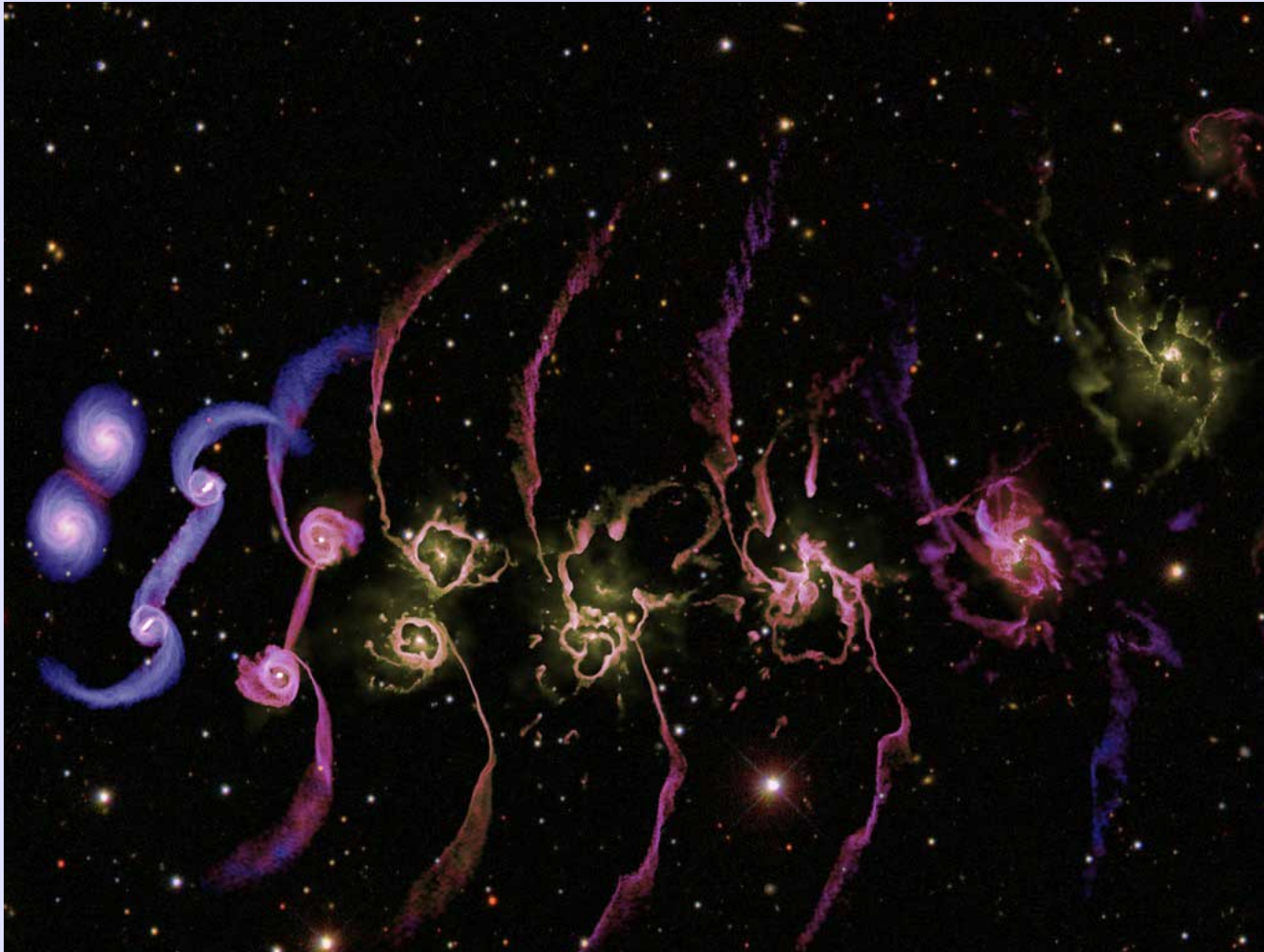


Super massive blackhole binaries

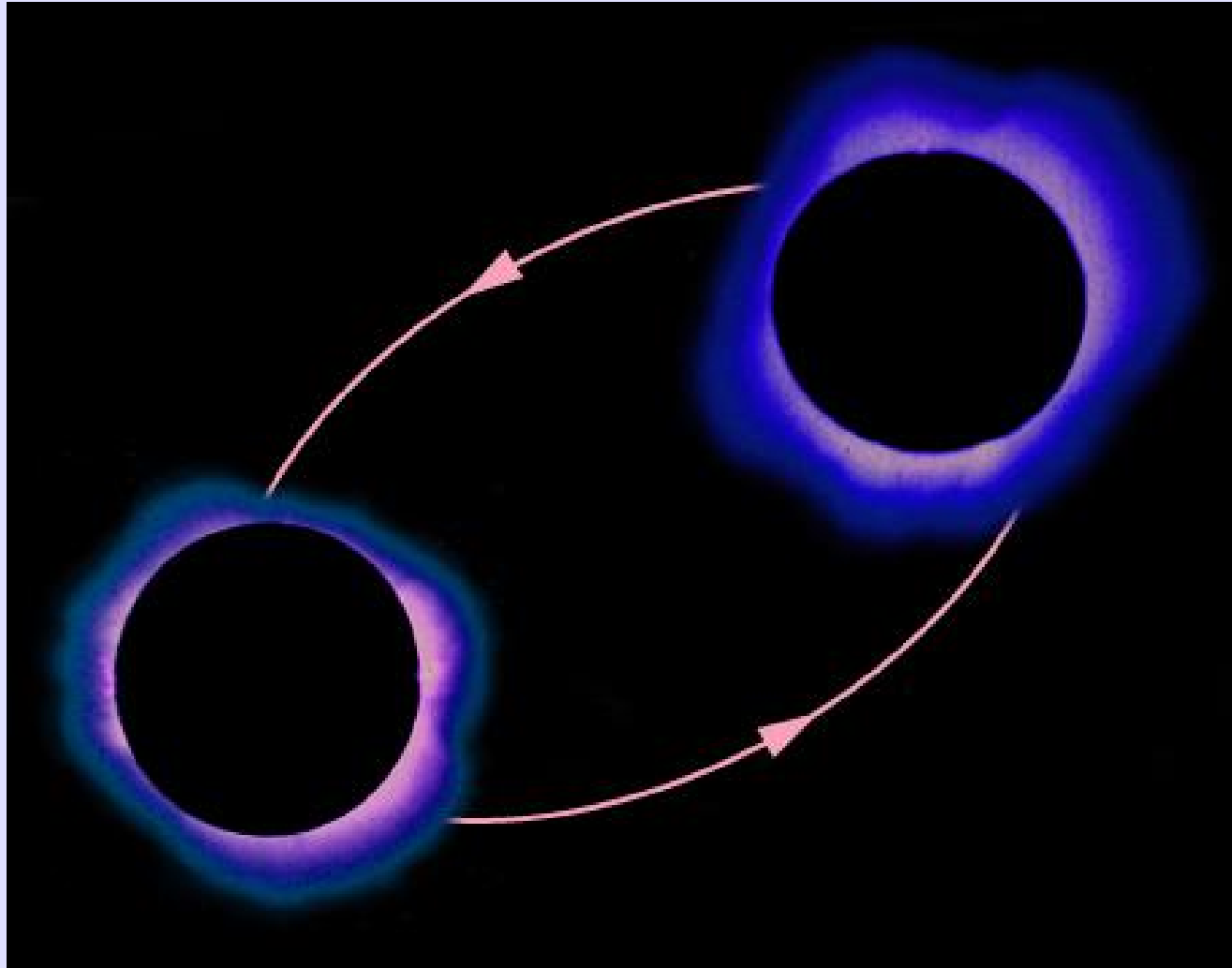
- ☞ There are strong evidence of existence of super massive black-hole of mass 10^5 to $10^8 M_{\odot}$ at the center of most of galaxies.

Super massive blackhole binaries

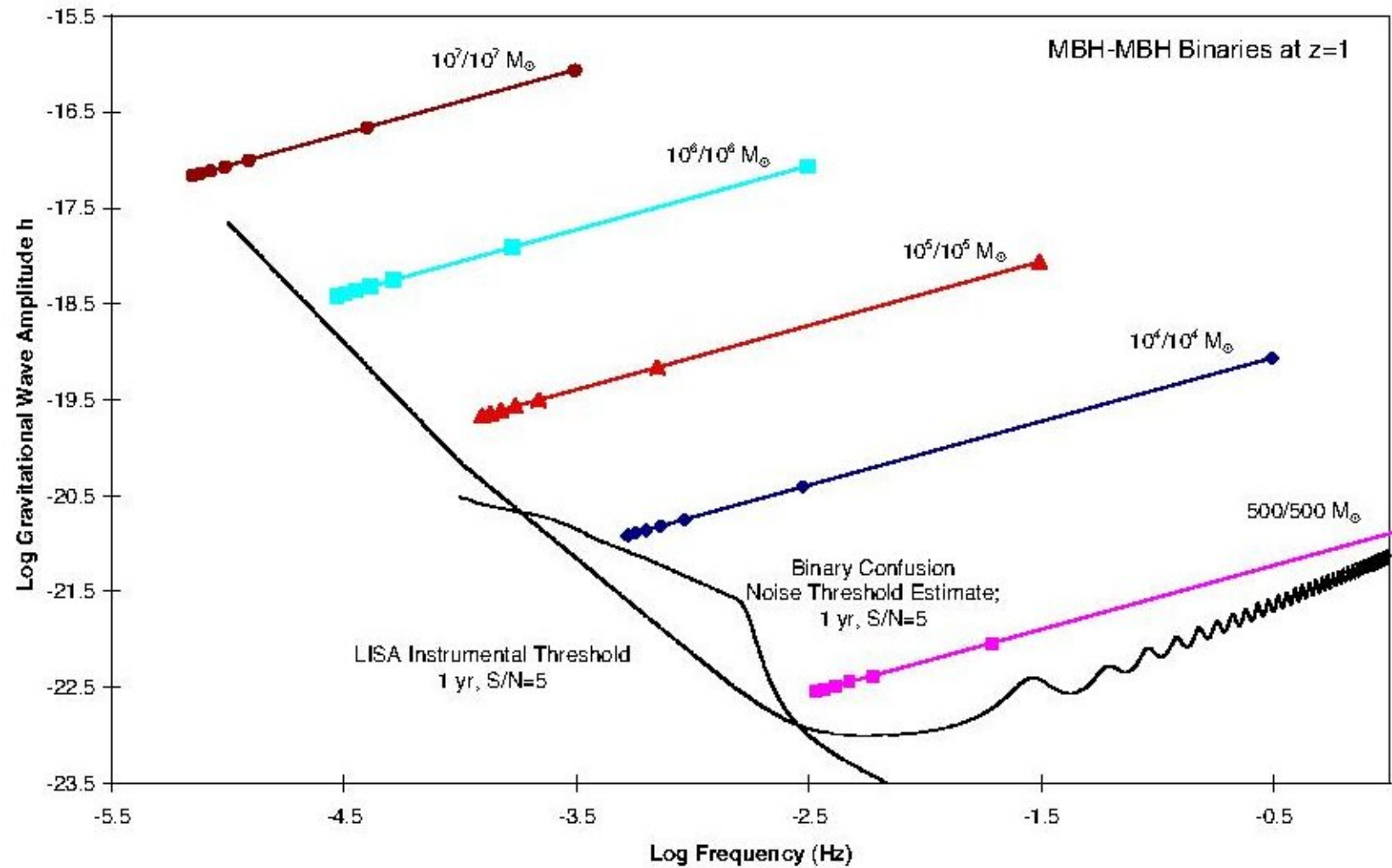
- ☞ There are strong evidence of existence of super massive black-hole of mass 10^5 to $10^8 M_{\odot}$ at the center of most of galaxies.
- ☞ It is seen that Galaxies collide,



Binary blackhole's are one of promising sources in the LISA band, can be seen almost anywhere in the Universe.



Sensitivity curve for blackhole binaries:



Super-massive blackhole binaries are very important sources in the LISA band.

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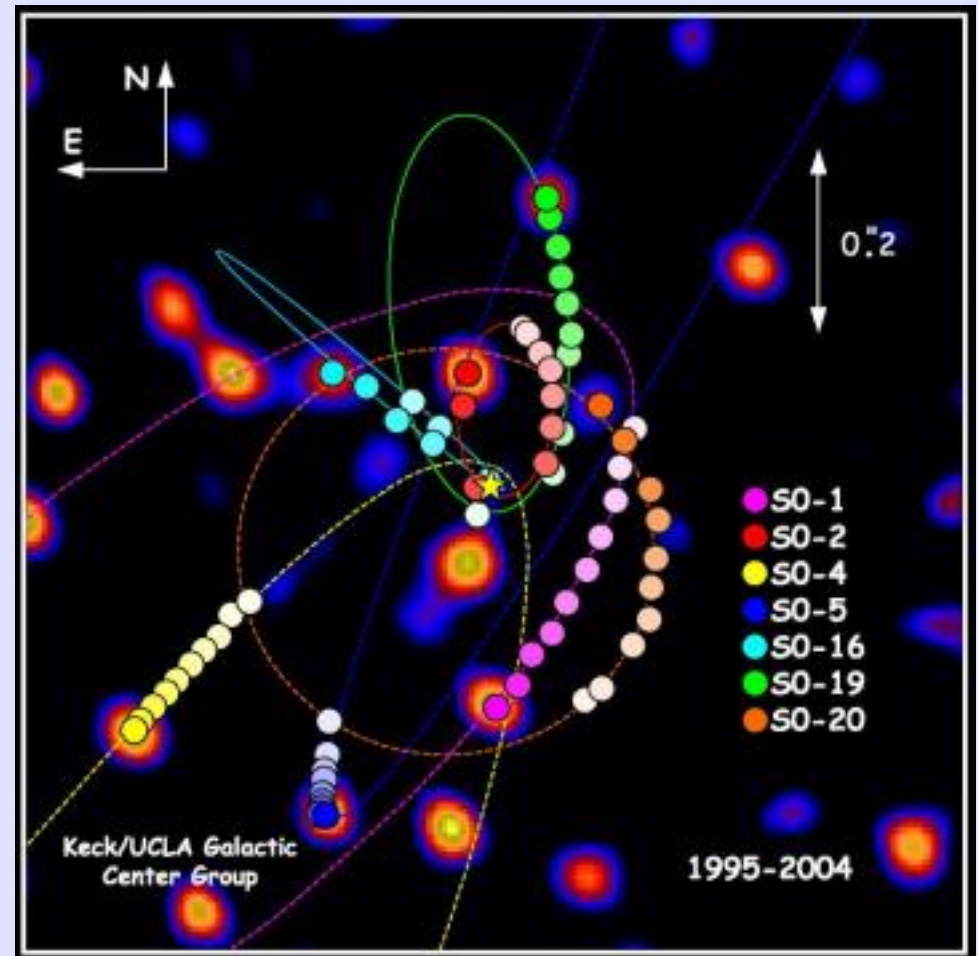
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- ☞ The binary blackholes can be used to make an independent estimate of Cosmological parameter.
- ☞ They serve as test for Einstein theory in the strong field limit.
- ☞ Can be used to test the alternative theory of gravity.

Extreme mass ratio inspiral(EMRI's)



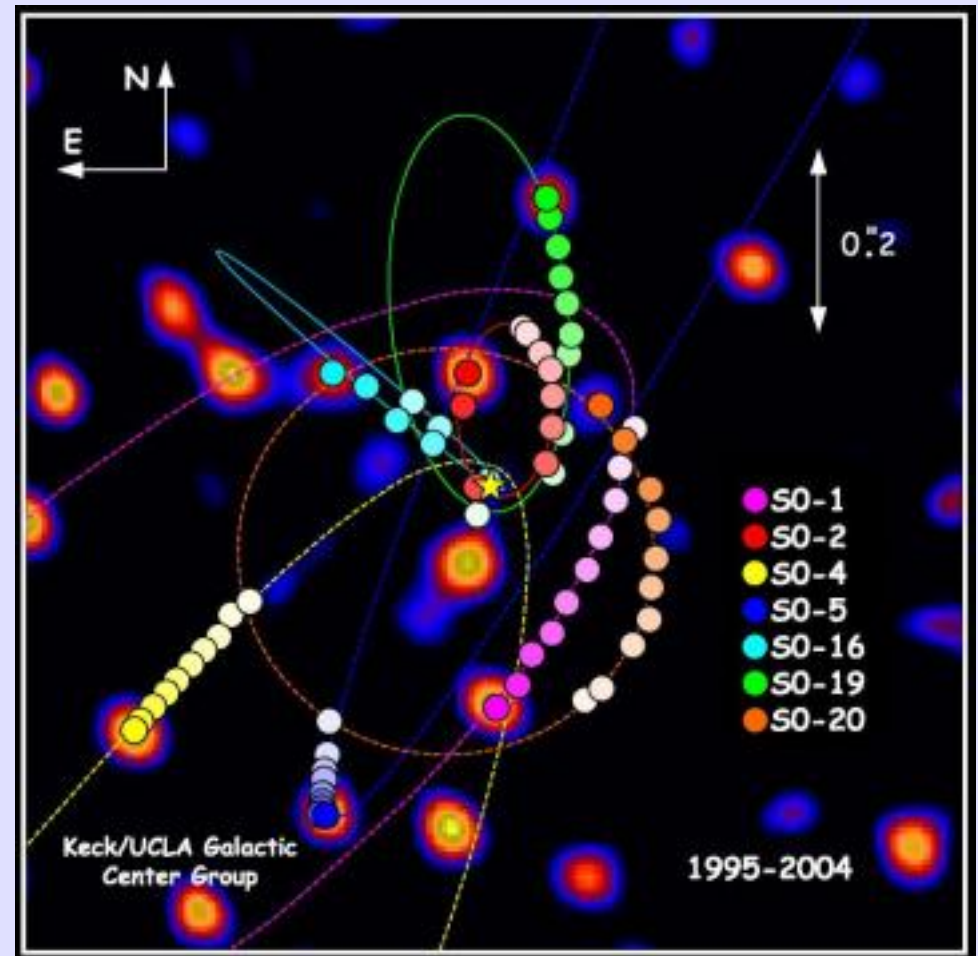
Galactic centers, such as our Galaxy hosts blackhole of mass $10^5 M_{\odot}$. Some of the stars orbiting around them can fall in to blackholes emitting gravitational waves:



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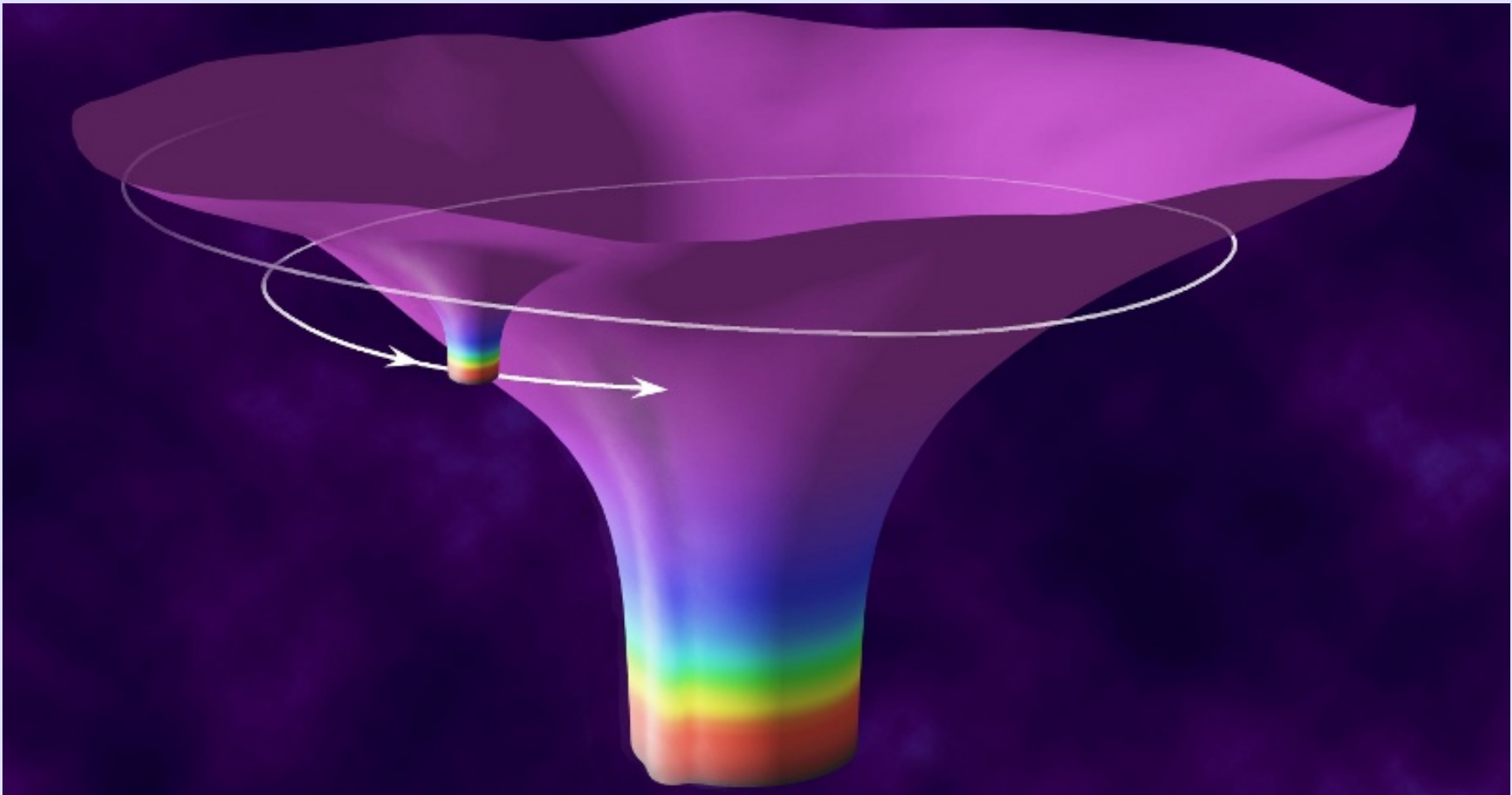


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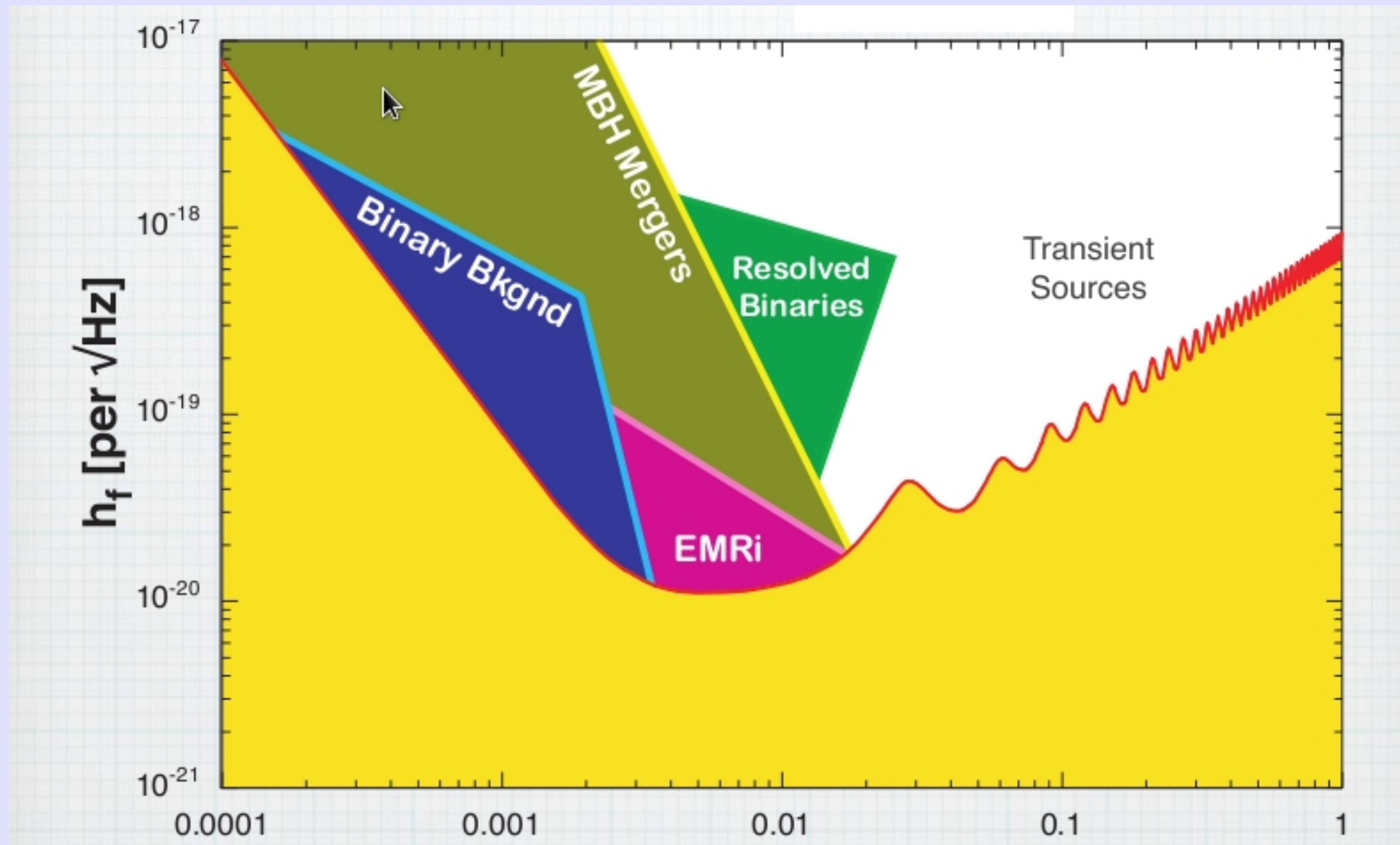


Or one can have blackholes of mass $1 - 10 M_{\odot}$ falling into central blackhole

- ☞ We have small objects of mass 1 to 50 M_{\odot} falling in to black-holes of mass $10^6 M_{\odot}$ is called EMRI



Sensitivity curve for EMRI:



In order to observe astrophysically important sources such as these, one needs a space mission such as LISA.

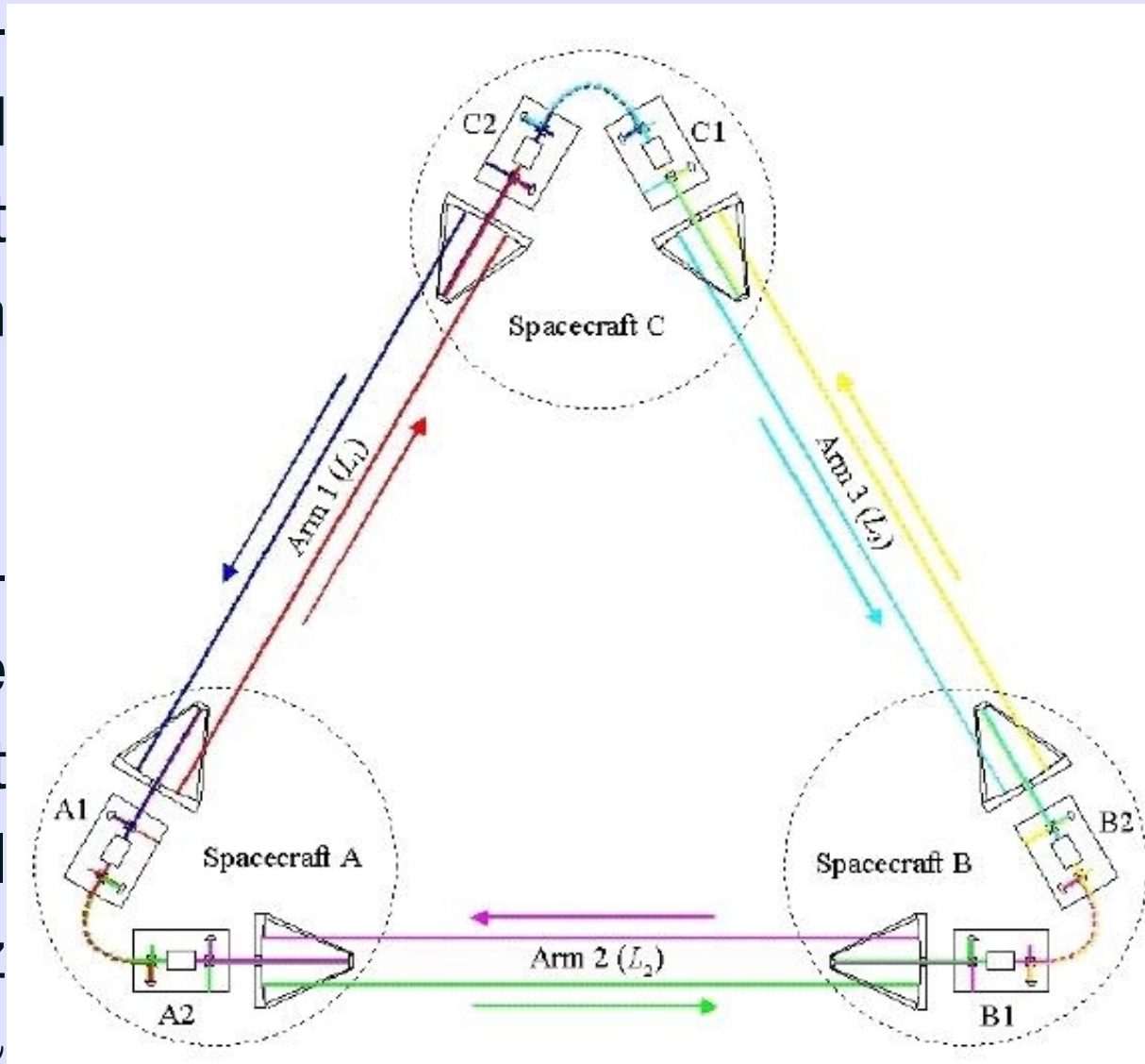
Mission overview

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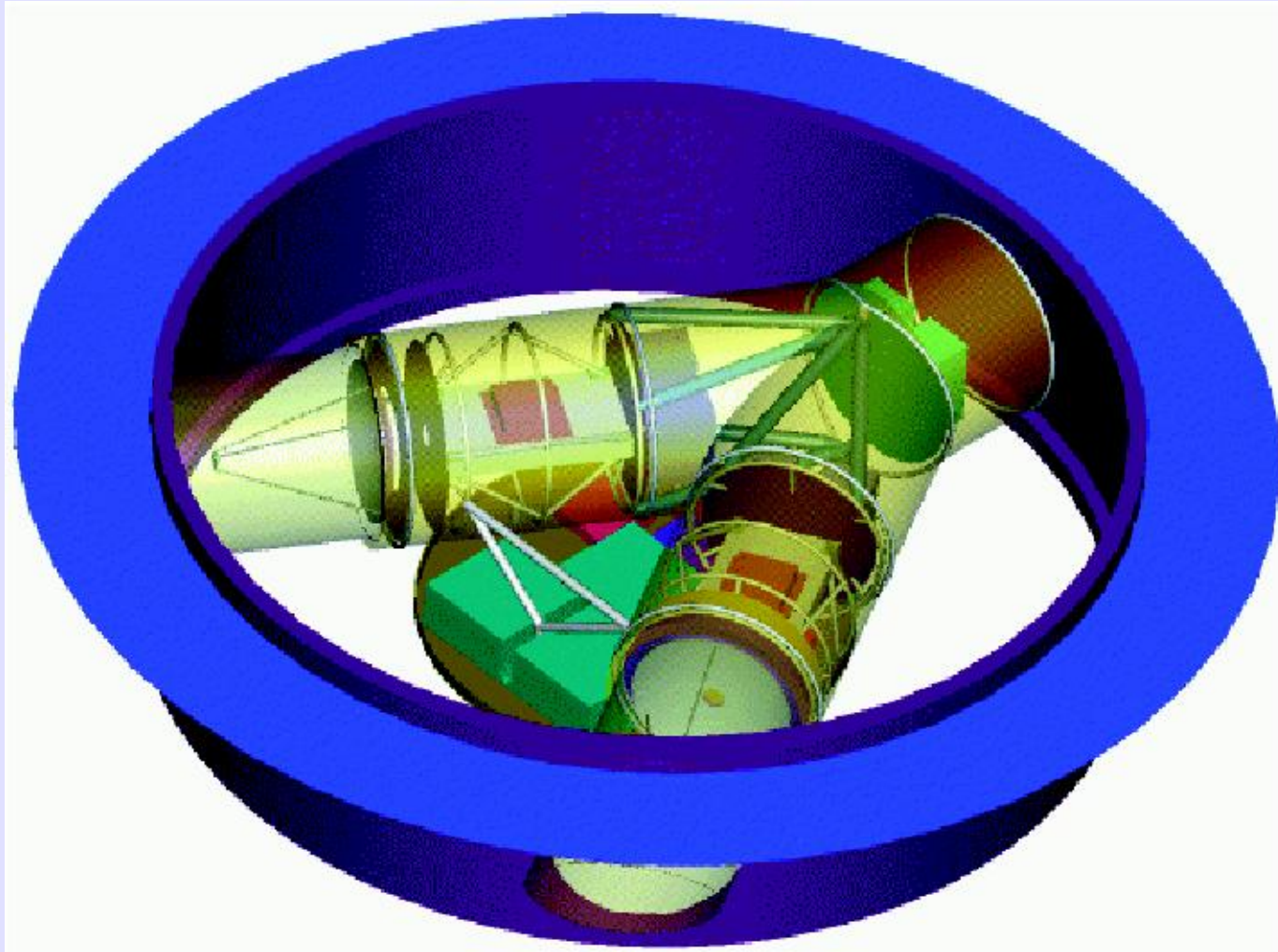
- ☞ The mission LISA consists of three identical spacecrafts orbiting around the Sun forming a giant laser interferometer of arm length 10^6 KM

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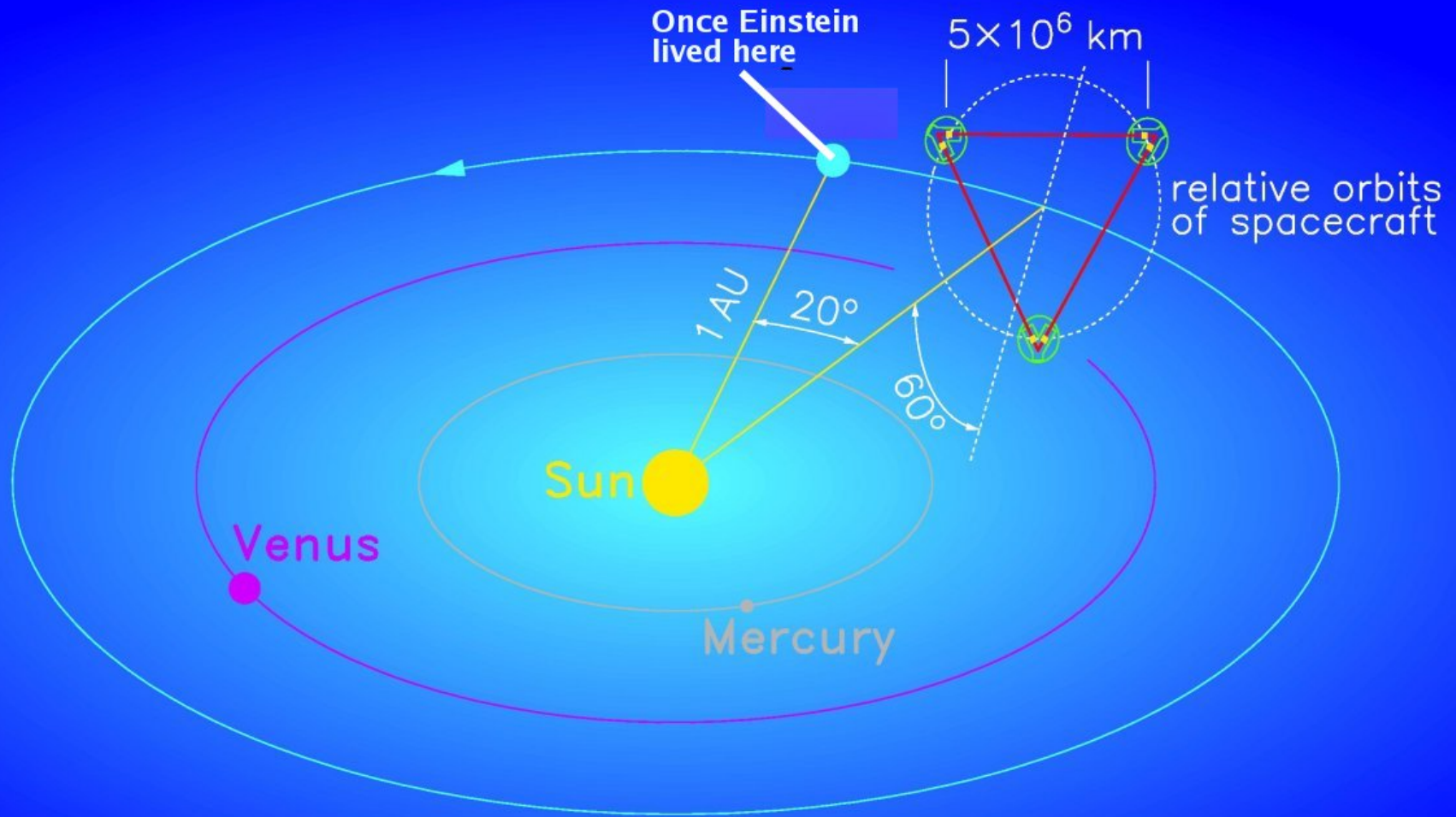
- ✎ The mission LISA consists of three identical spacecraft orbiting around the Sun forming a giant laser interferometer of arm length 10^6 KM
- ✎ Six laser beams are exchanged between these three spacecraft to detect gravitational wave signal in the band 10^{-4} - 1 Hz with peak sensitivity of $h \sim 10^{-23}$.



Each spacecraft hosts, two test masses and optical system to communicate with the distance spacecraft



LISA orbits



LISA orbits

- ➡ Each of the LISA spacecraft go around Sun, in elliptical orbits such that entire formation remain almost equilateral triangle.
- ➡ The plane of the LISA triangle makes an angle of 60° with the ecliptic plane.
- ➡ The center of this LISA constellation moves around the Sun in an earth-like orbit ($R = 1\text{AU}$), 20° behind the Earth.
- ➡ The triangular constellation revolves once around its center as it completes one orbit in the course of one year.

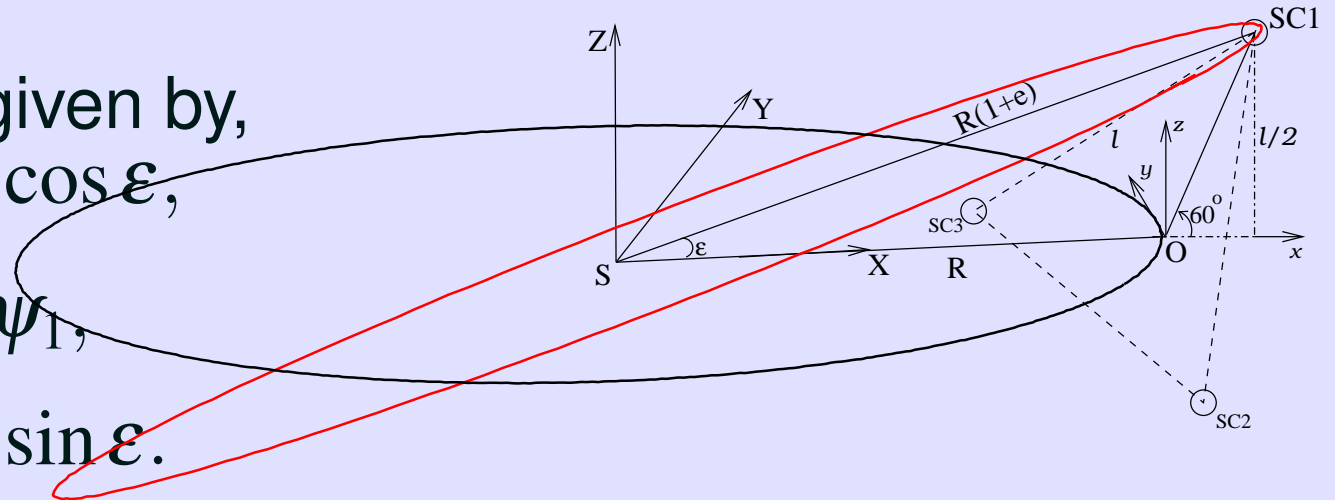
A choice of orbits

A choice of orbits are given by,

$$X_1 = R(\cos \psi_1 + e) \cos \epsilon,$$

$$Y_1 = R\sqrt{1 - e^2} \sin \psi_1,$$

$$Z_1 = R(\cos \psi_1 + e) \sin \epsilon.$$



$$\tan \epsilon = \frac{\alpha}{1 + \alpha/\sqrt{3}}, \quad \alpha = \frac{l}{2R}$$

$$e = \left(1 + \frac{2}{\sqrt{3}}\alpha + \frac{4}{3}\alpha^2 \right)^{1/2} - 1$$

$$\psi_1 + e \sin \psi_1 = \Omega t$$

Orbits for other two spacecraft can be obtained by rotating this by 180° about z -axis

CW-Frame and Equations

Clohessy and Wiltshire or Hill's equations are linearised dynamical equations for test-particles in the neighborhood of reference point, in our case the LISA centroid. These equations are written in a frame which has its origin on the reference orbit and also rotates with angular velocity Ω . The equation for a free test particle are given by,

$$\ddot{x} - 2\Omega\dot{y} - 3\Omega^2x = 0,$$

$$\ddot{y} + 2\Omega\dot{x} = 0,$$

$$\ddot{z} + \Omega^2z = 0.$$

The solutions:

$$x(t) = \frac{\dot{x}_0}{\Omega} \sin \Omega t - \left(3x_0 + \frac{2\dot{y}_0}{\Omega} \right) \cos \Omega t + 2 \left(2x_0 + \frac{\dot{y}_0}{\Omega} \right)$$

$$y(t) = \left(6x_0 + \frac{4\dot{y}_0}{\Omega} \right) \sin \Omega t + \frac{2\dot{x}_0}{\Omega} \cos \Omega t - 3(2\Omega x_0 + \dot{y}_0)t \\ + \left(y_0 - \frac{2\dot{x}_0}{\Omega} \right)$$

$$z(t) = z_0 \cos \Omega t + \frac{\dot{z}_0}{\Omega} \sin \Omega t$$

Ignoring runaway solutions and offset solution the condition for stable configuration is given by

$$z_0 = \mu \sqrt{3} x_0 \quad \text{and} \quad \frac{\dot{z}_0}{\Omega} = \frac{1}{2} \mu \sqrt{3} y_0, \quad \mu = \pm 1$$

The solutions are given by,

$$x(t) = \frac{1}{2}\rho_0 \cos(\Omega t - \phi_0),$$

$$y(t) = -\rho_0 \sin(\Omega t - \phi_0),$$

$$z(t) = \mu\rho_0 \frac{\sqrt{3}}{2} \cos(\Omega t - \phi_0),$$

where

$$\rho_0 = \sqrt{4x_0^2 + y_0^2} \quad \tan \phi_0 = \frac{y_0}{2x_0}$$

The orbits given earlier when approximated to first order in $\alpha = l/(2R)$ and transformed to CW frame we get:

$$x_k = eR \cos \left[\Omega t - (k-1) \frac{2\pi}{3} \right]$$

$$y_k = -2eR \sin \left[\Omega t - (k-1) \frac{2\pi}{3} \right]$$

$$z_k = \sqrt{3}eR \cos \left[\Omega t - (k-1) \frac{2\pi}{3} \right]$$

we identify $\rho_0 = 2eR$ and $\phi_0 = 2\pi(k-1)/3$

The general result is

In the CW frame there are just two planes which make angles of $\pm\pi/3$ with the (x-y) plane, in which test particles obeying CW equations perform rigid rotations about the origin with angular velocity $-\Omega$.

S. V. Dhurandhar et. al, “Fundamentals of the LISA Stable Flight Formation”, *Class. Quantum Grav.* **22** 481(2005).

Time-Delay interferometer

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- ❏ Because different lasers are used at eand points, laser noise play important role.

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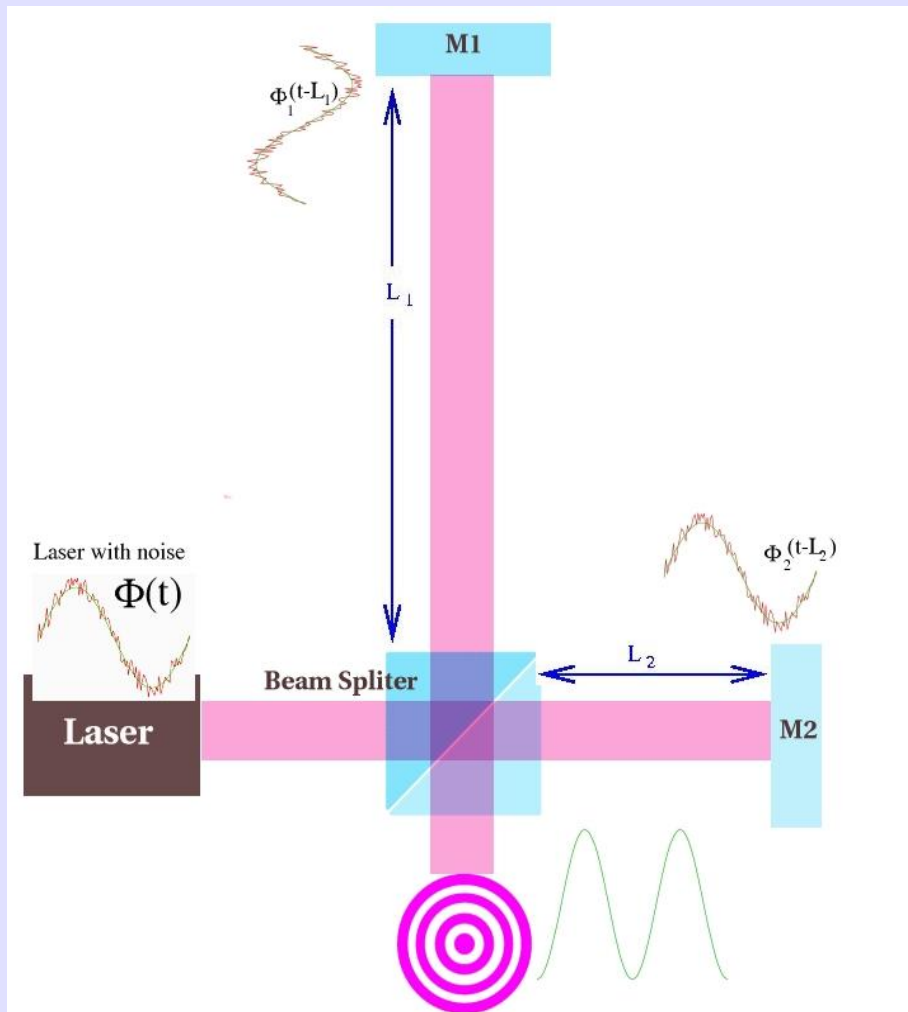
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- ☞ In the case of ground based detectors, where both arms are exactly of same length (integral multiple of $\frac{\lambda}{2}$) and the laser noise cancels out.
- ☞ Time-Delay interferometry is one of novel technique suggested by Armstrong et. al (J.W. Armstrong, F.B Estabrook and M. Tinto, *Astrophys. J.* **527**, 814(1999).) for constructing unequal arm interferometer in space
- ☞ In this method the individual beams are combined off-line after introducing suitable time delay corresponding to the light travel time across the arms to simulate the interferometer.

Unequal Arm Interferometer

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Let $\Phi(t)$ be the Phase fluctuation of laser



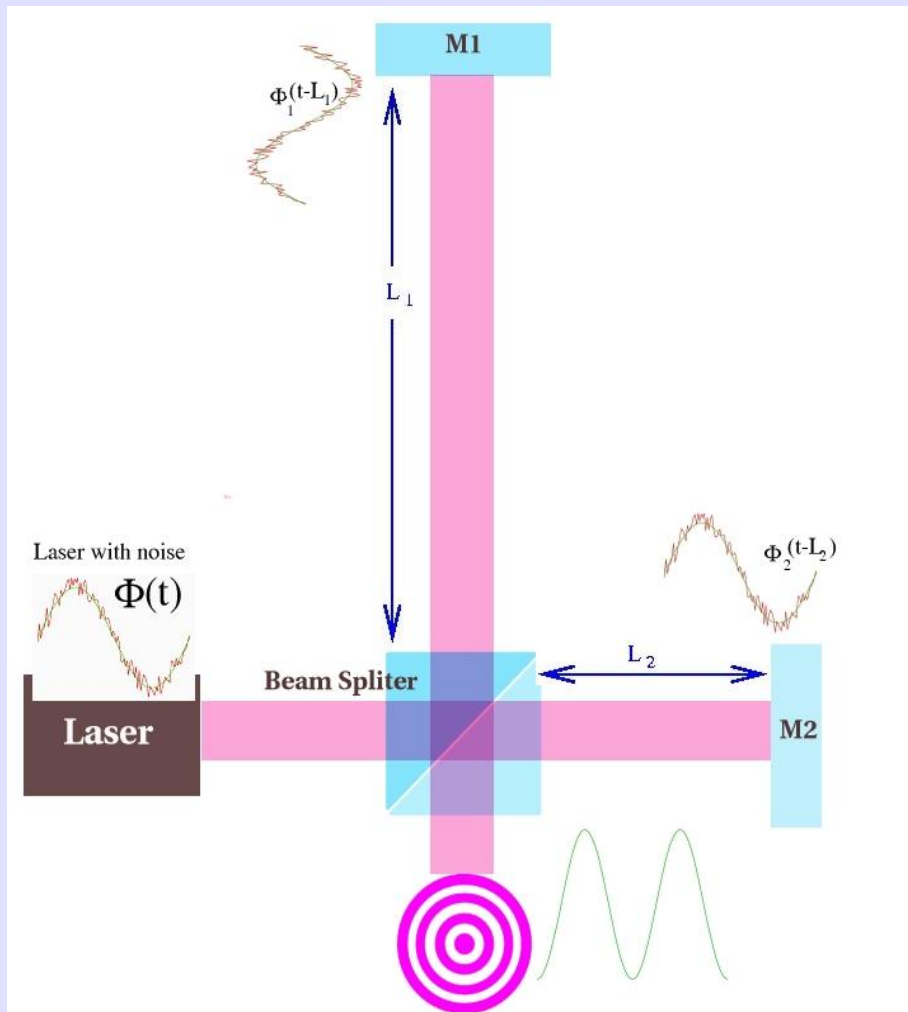
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If $L_1 \neq L_2$

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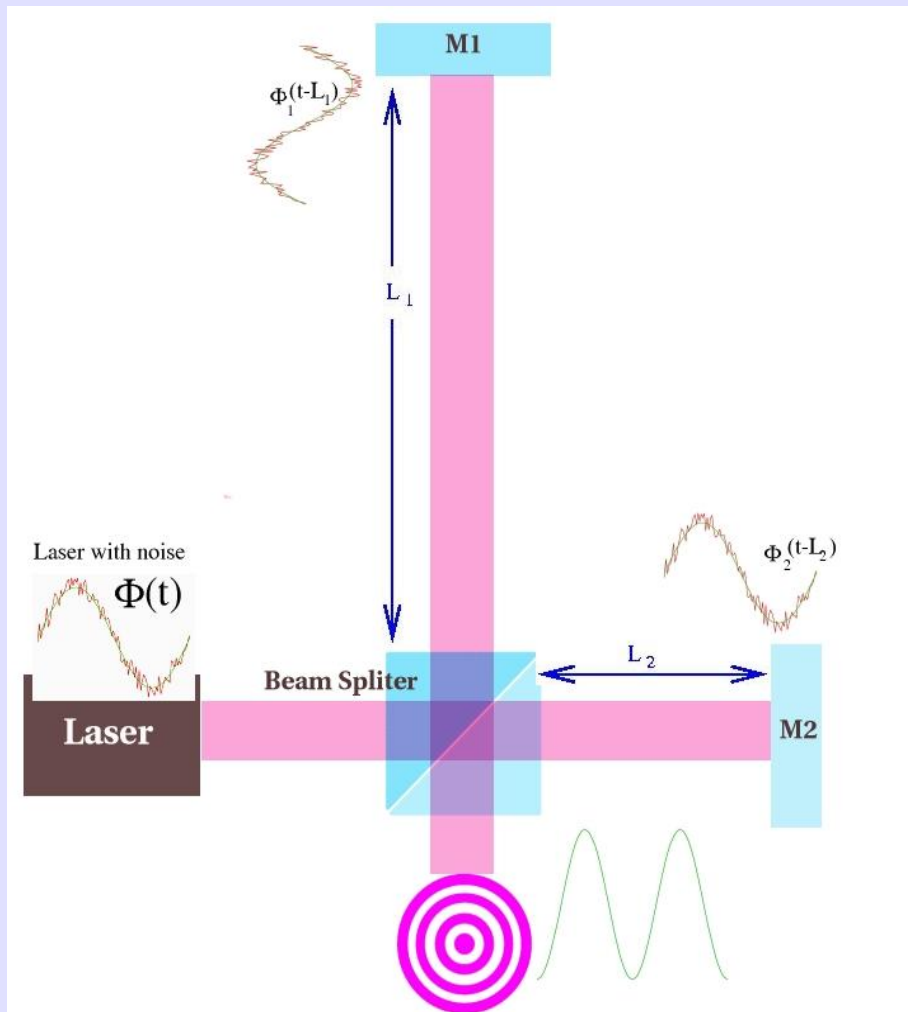
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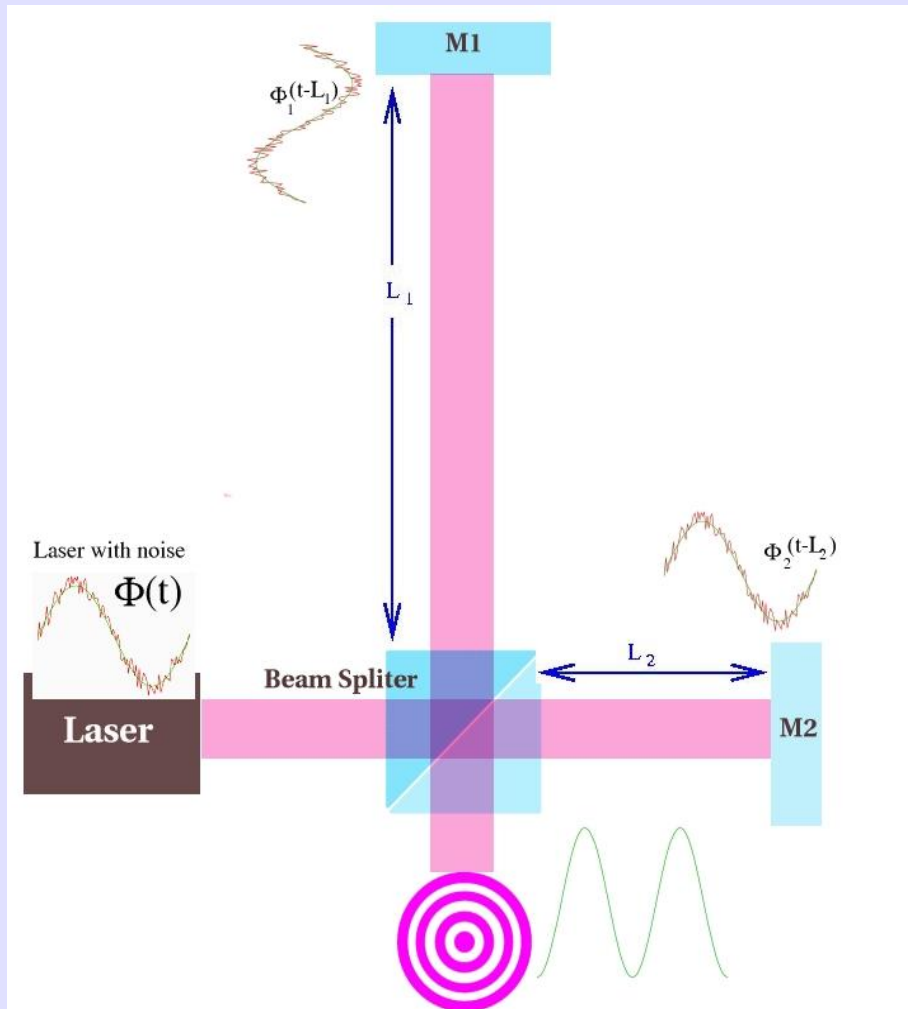
$$\Phi_2(t) = \Phi(t - 2L_2) - \Phi(t)$$

Clearly,

$$\Phi_1(t) - \Phi_2(t) \neq 0$$



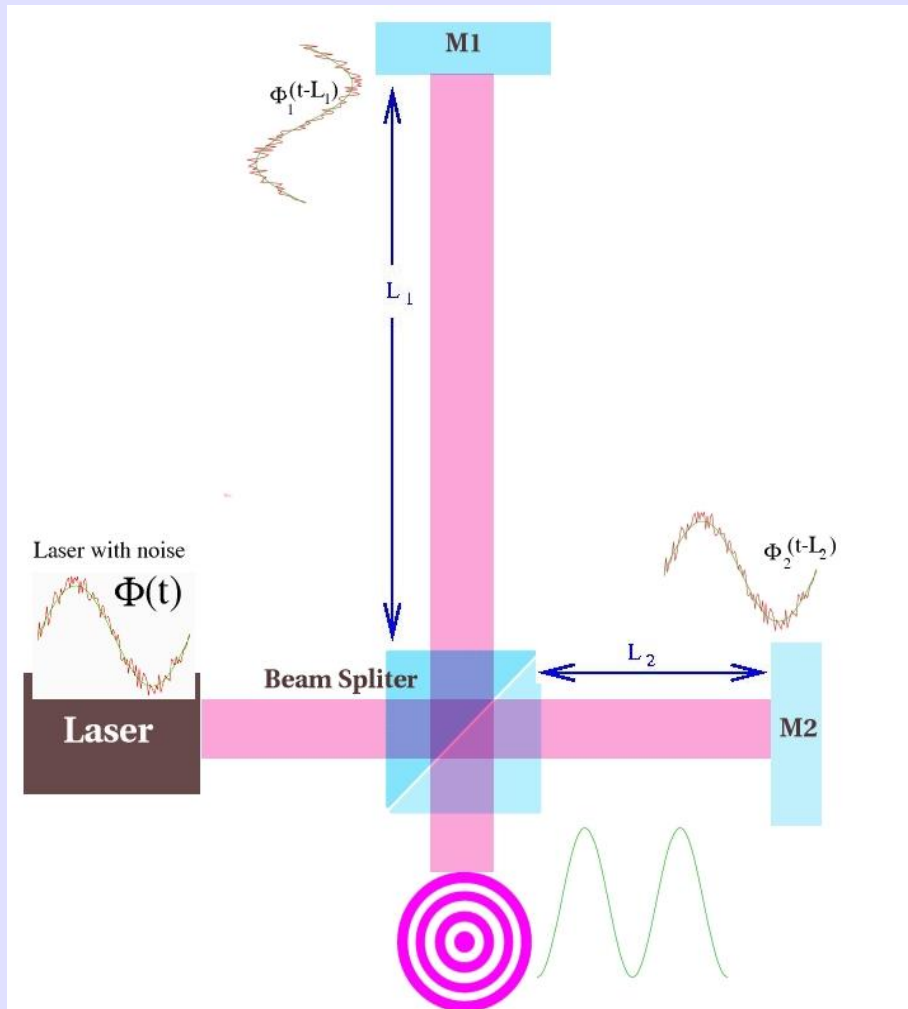
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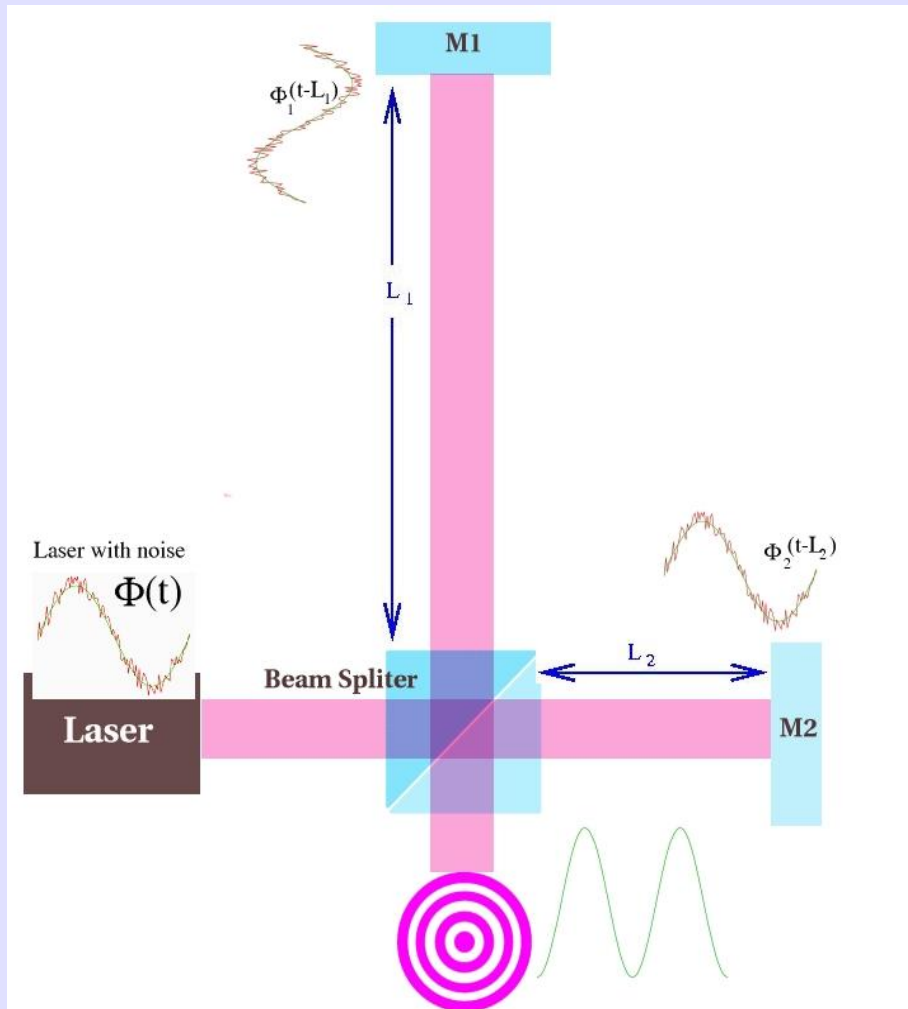
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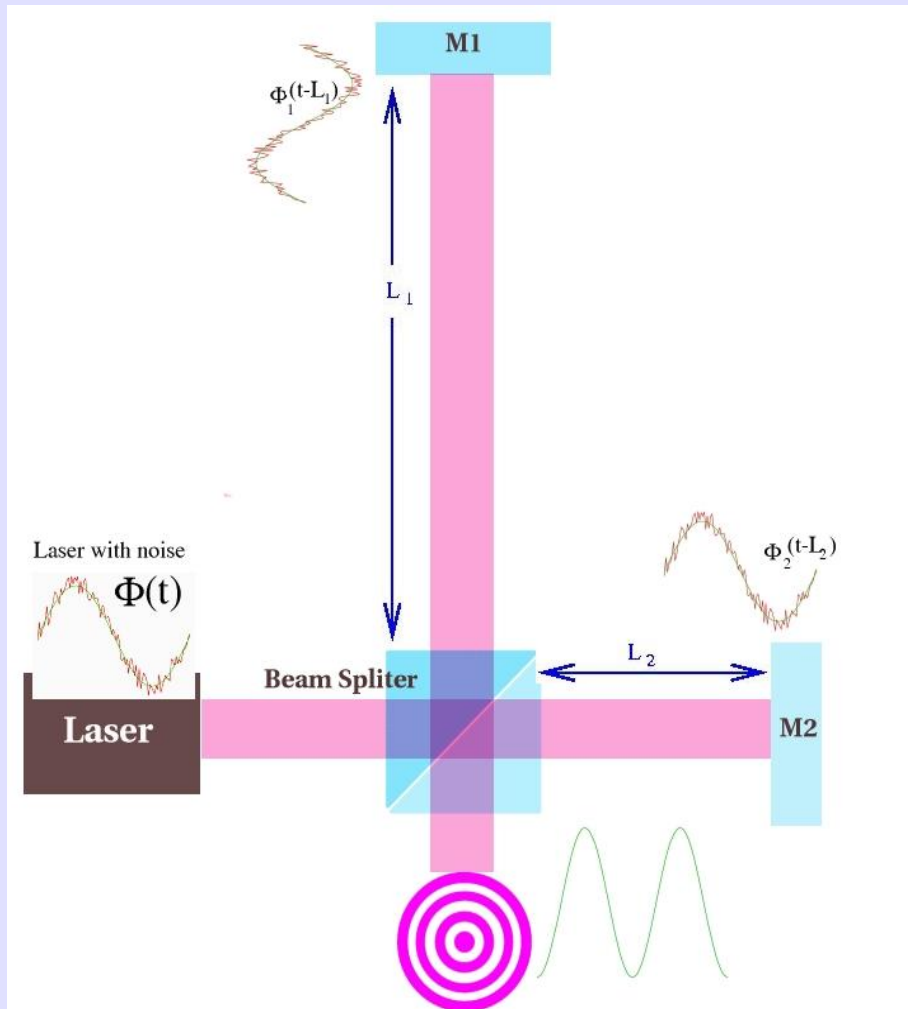
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LISA is a more complex system. There are three arms and six beams.



Delay Operator

S. V. Dhurandhar et. al, “Algebraic approach to time-delay data analysis for LISA”, *Phs Rev. D*65, 102002(2002), gr-qc/0112059.

Let $a(t)$ be any arbitrary function of time and L_k be the length of the k th arm, then we define time delay operator for the arm k as,

$$D_k a(t) = a(t - L_k)$$

$c = 1$ and all the distances are measured in the unit of time.

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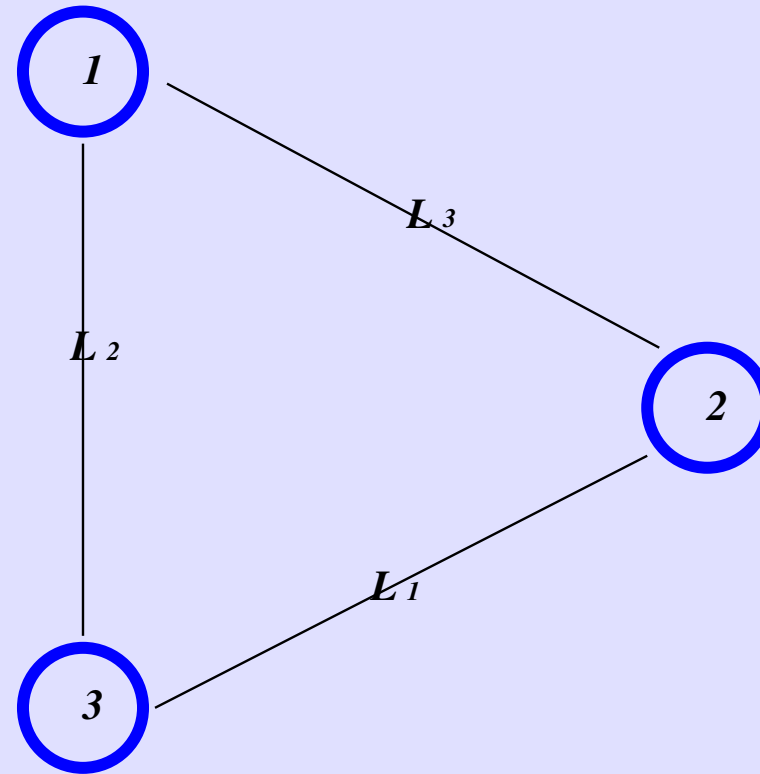
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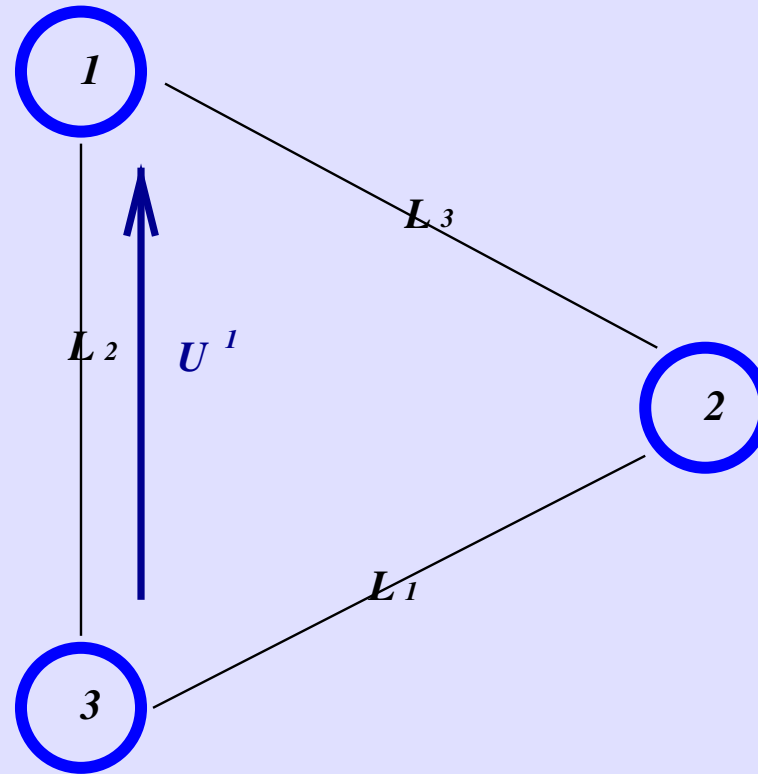
They commute: $D_1^k D_2^l = D_2^l D_1^k$

Back to LISA geometry and beams



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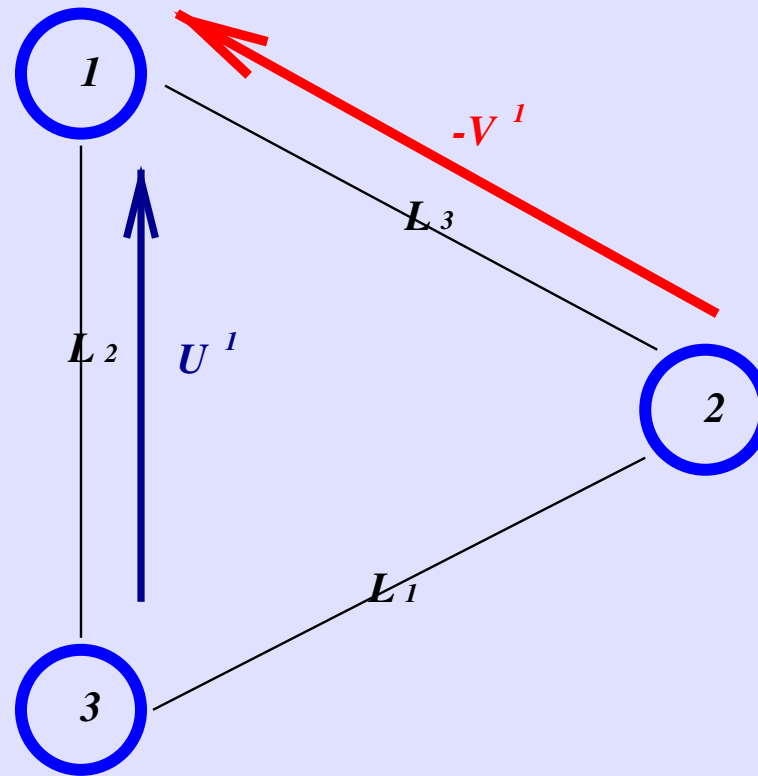
$$U^1 = D_2 C_3 - C_1$$



Back to LISA geometry and beams

$$U^1 = D_2 C_3 - C_1$$

$$V^1 = C_1 - D_3 C_2$$



Back to LISA geometry and beams

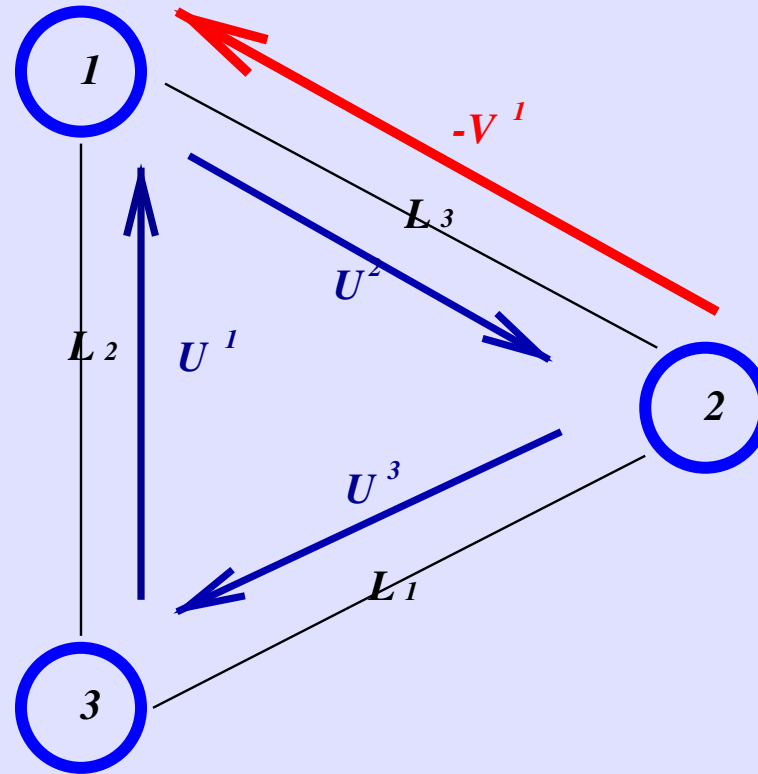
U_i's in +ve direction

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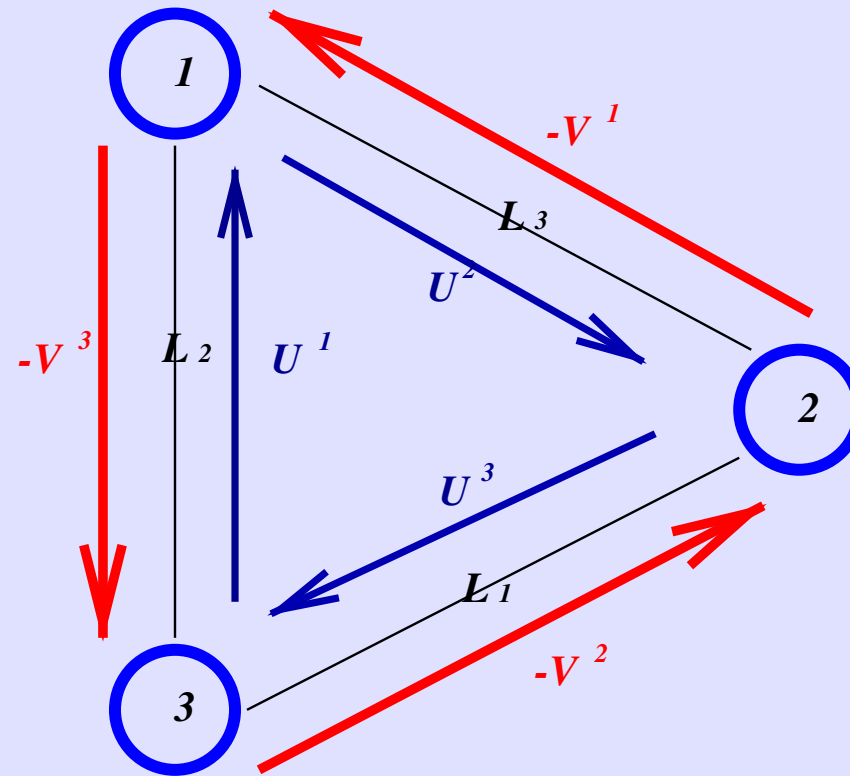
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V_i 's in -ve direction

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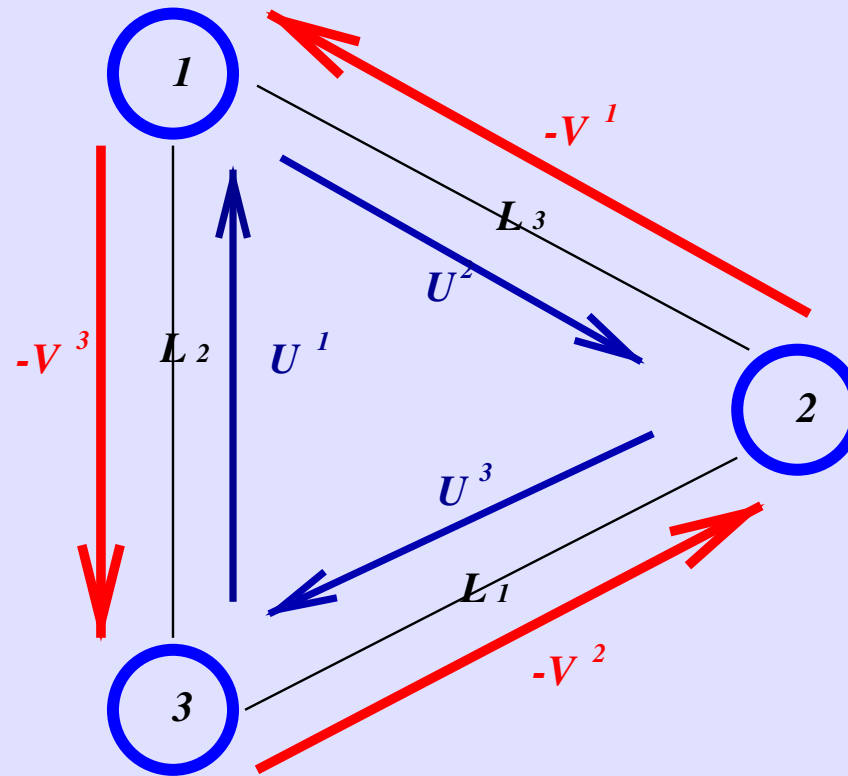
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There are six beams



Redundancy in Data

Back to LISA geometry and beams

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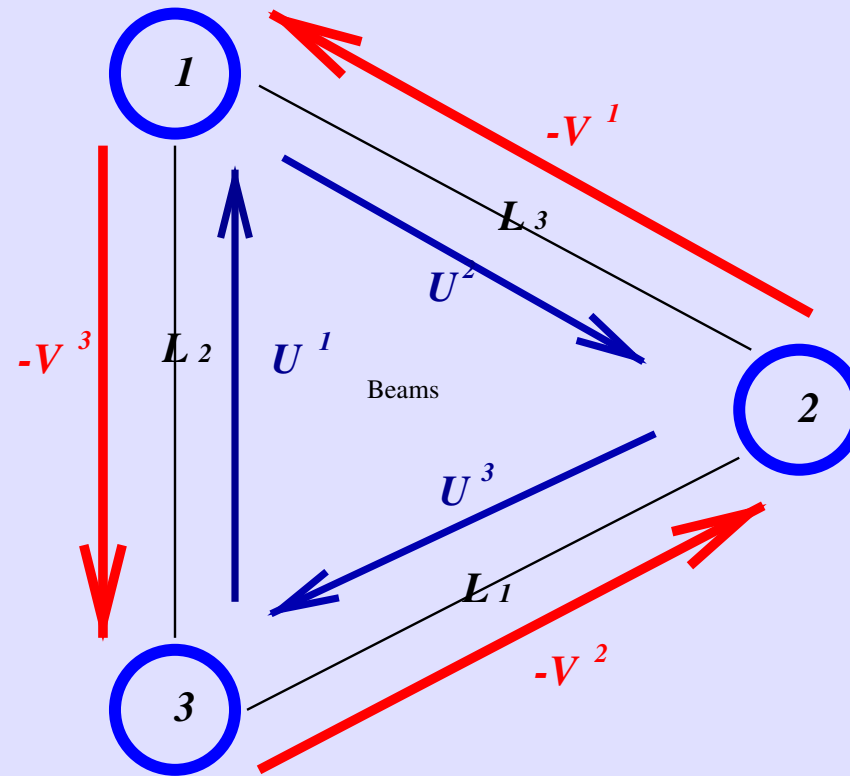
$$U^3 = D_1 C_2 - C_3$$

V_i's in -ve direction

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$$V^3 = C_3 - D_2 C_1$$

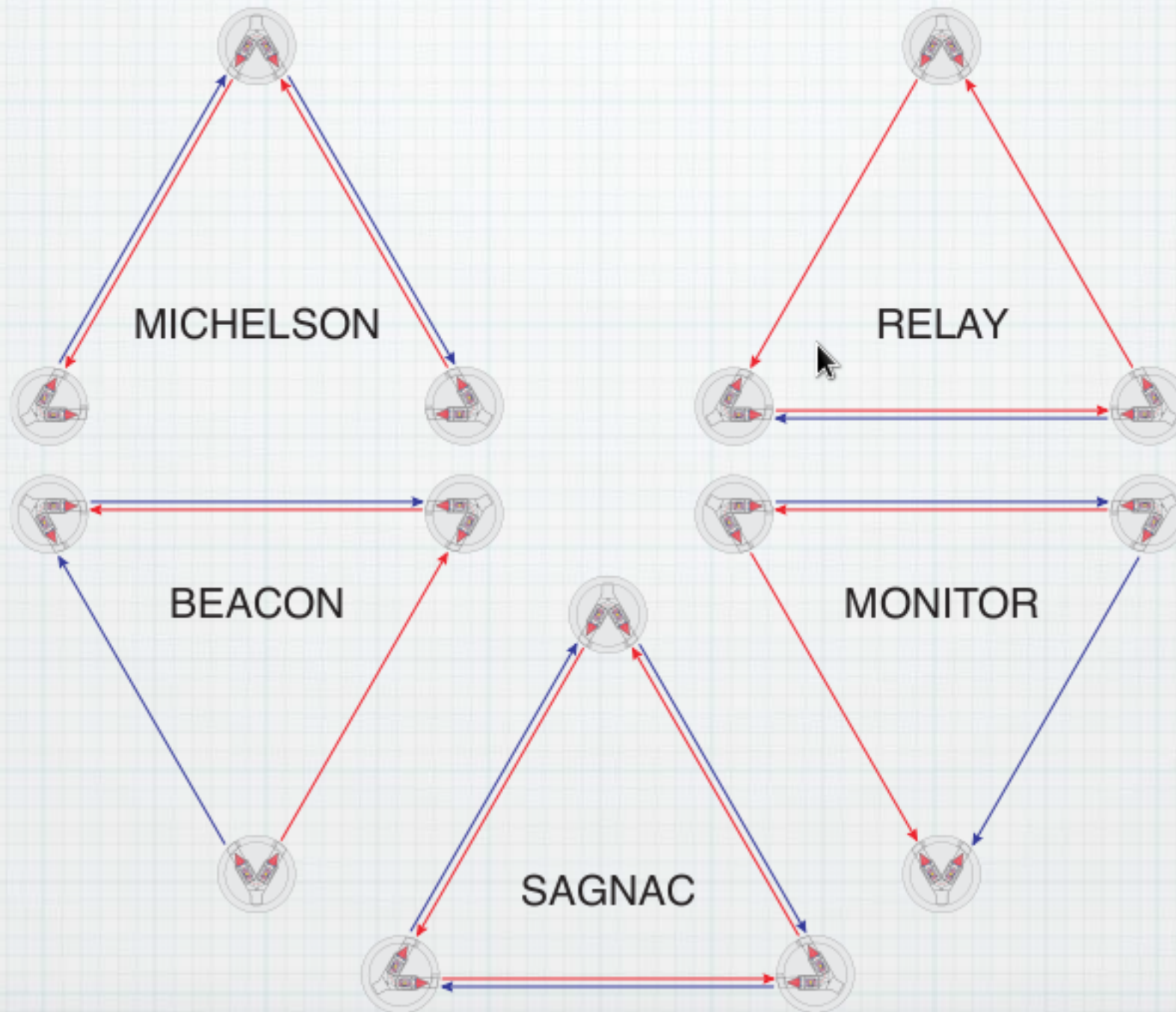


There are six beams



Redundancy in Data

Known Noise cancellation solution



Laser noise Cancellation Data combination

A general data combination is given by the combination of U^i and V^i of the form

$$X = \sum_{i=1}^3 p_i V^i + q_i U^i$$

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A noise cancellation Data combination, we need to determine p_i and q_i such that,

$$\sum_{i=0}^3 p_i V^i + q_i U^i = 0$$

We need to solve for (p_i, q_i) as functions of D_i .

The solution to this equation is well known in the algebra and forms a module called “***First Module of Syzygies***” , over the polynomial ring D_i .

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They are called generators! and any element of module can be written of the form $\sum p_i g^i$ where $p_i \in \mathcal{R}$ and g^i 's are the generators.

The Generators

The solutions are represented by generator with (p_i, q_i) ,

$$X^{(1)} = \alpha = (1, D_3, D_1 D_3, 1, D_1 D_2, D_2),$$

$$X^{(2)} = \beta = (D_1 D_2, 1, D_1, D_3, 1, D_2 D_3),$$

$$X^{(3)} = \gamma = (D_2, D_2 D_3, 1, D_1 D_3, D_1, 1),$$

$$X^{(4)} = \zeta = (D_1, D_2, D_3, D_1, D_2, D_3).$$

With these generator any solution can be expressed as:

$$X(p_i, q_i) = \sum_{I=1}^4 \alpha_{(I)} X^{(I)}$$

In general, $\alpha_{(I)}$ can be polynomials in D_i

Advantages of the Formalism

- ☞ All the Noise cancellation combination can be obtained using a set of generators. Any noise cancellation data combination can be written as,

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